

Mass Flux in vertical
direction: N_z

$$N_z = -WC - \langle wc \rangle - D \partial C / \partial z$$

convection by mean velocity convection by fluctuating velocity molecular diffusion

where $(\tilde{W}, \tilde{C}) = (W, C) + (w, c)$
instantaneous mean fluctuation

W, w : vertical velocity
 D : molecular diffusion coefficient
 C, c : concentration
 z : distance in upward direction
 $\langle \rangle$: time average

$D \partial C / \partial z$: negligible

→ N_z is independent of
diffusion coefficient

At the tropopause

$$N_z \sim -WC$$

Near the surface $W=0$

$$N_z \sim -\langle wc \rangle$$

Momentum flux τ_z :

$$\tau_z / \rho = -\langle uw \rangle \equiv (u^*)^2$$

Heat Flux q_z :

$$q_z / \rho c_p = -\langle w\theta \rangle \equiv u^* T^*$$

Mass flux N_z :

$$N_z = -\langle wc \rangle \equiv u^* C^*$$

Methods for flux measurement

1. Gradient method (profile method)

Boussinesq concept of

eddy diffusivity K_d

$$N_z = -\langle w c \rangle = K_d \partial C / \partial z = u^* C^*$$

$$-\langle w \theta \rangle = K_d \partial \theta / \partial z = u^* T^*$$

$$-\langle u w \rangle = K_m \partial U / \partial z = (u^*)^2$$

Near the surface level

1. u^* , T^* , C^* ; constant

2. $K_d = 1.2 \kappa u^* z / \phi_h$

$$K_m = \kappa u^* z / \phi_m$$

κ : Kármán constant (=0.40)

Shear functions, ϕ_m , ϕ_h

$$\phi_m \equiv (\kappa z / u^*) dU / dz,$$

$$\phi_h \equiv (\kappa z / T^*) d\theta / dz$$

are universal function

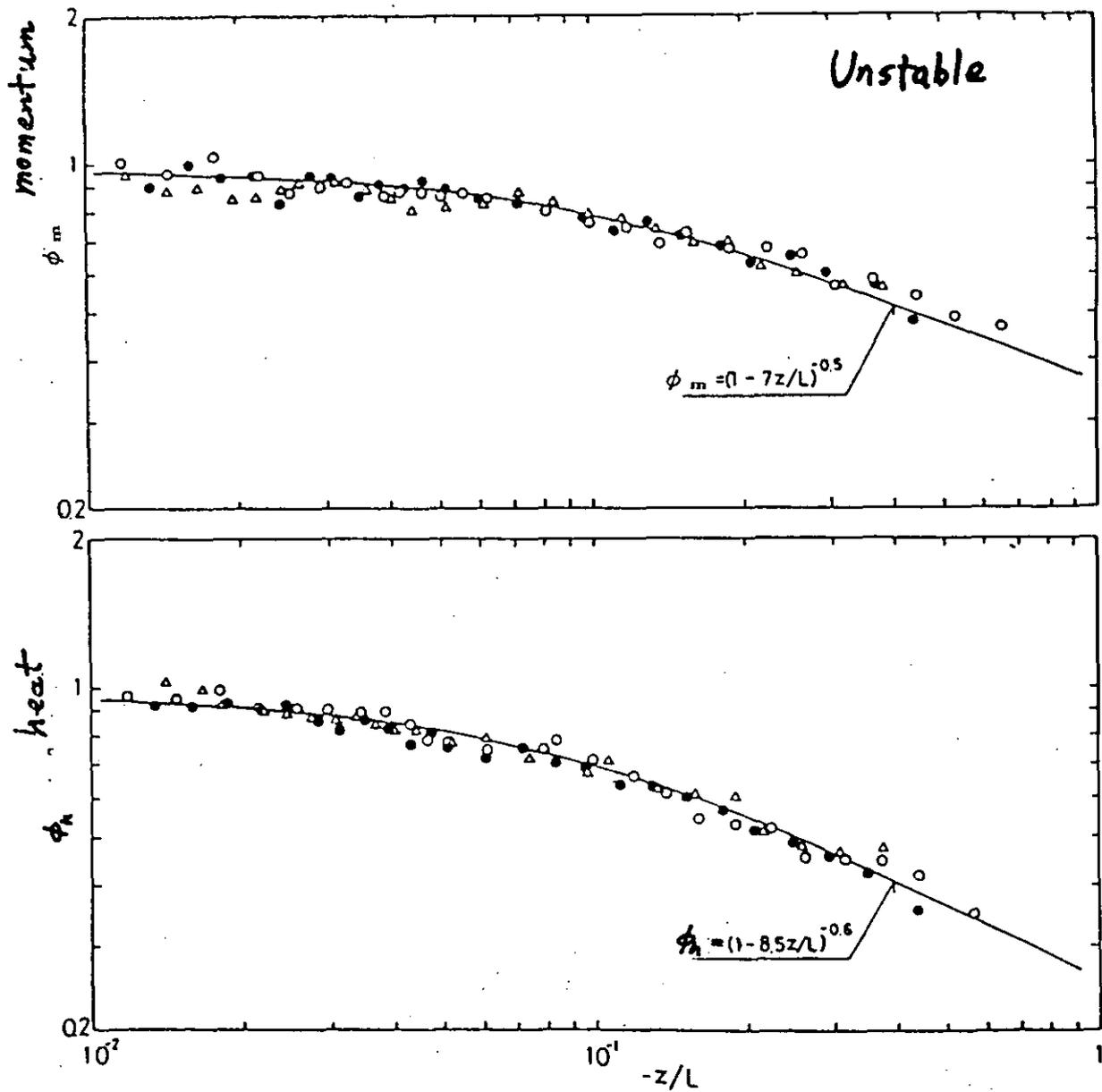
of z/L . stratification level, atmospheric stability

L : Monin-Obukhov length

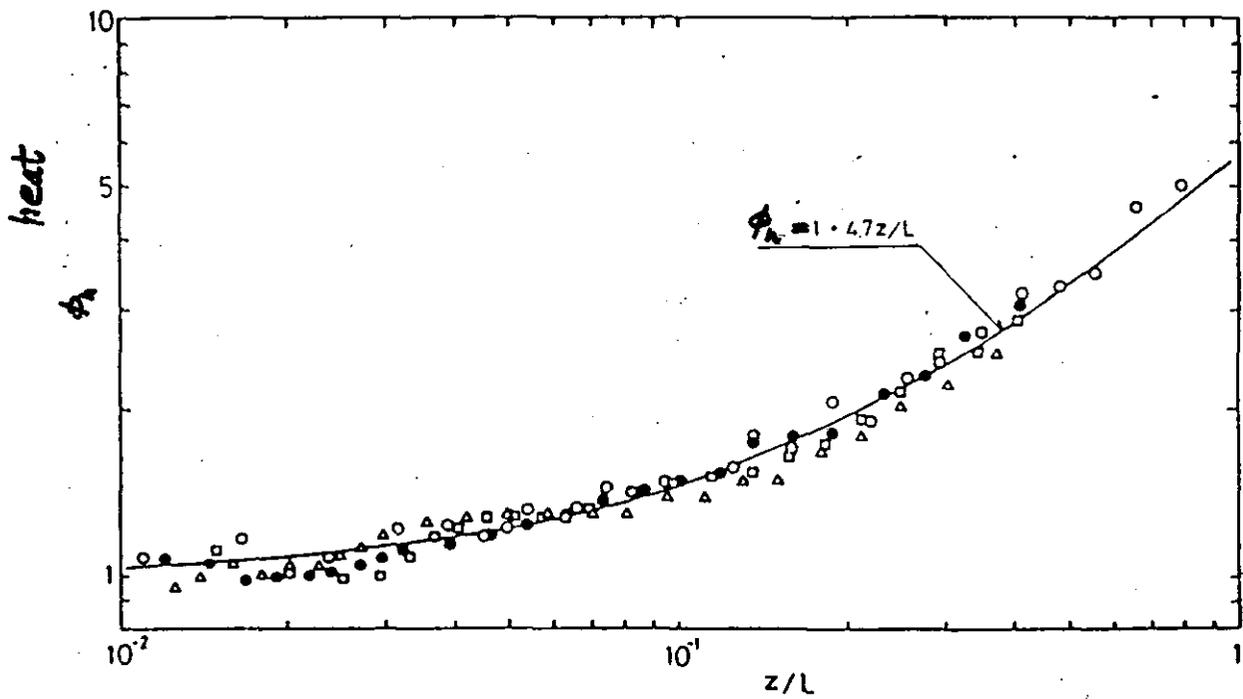
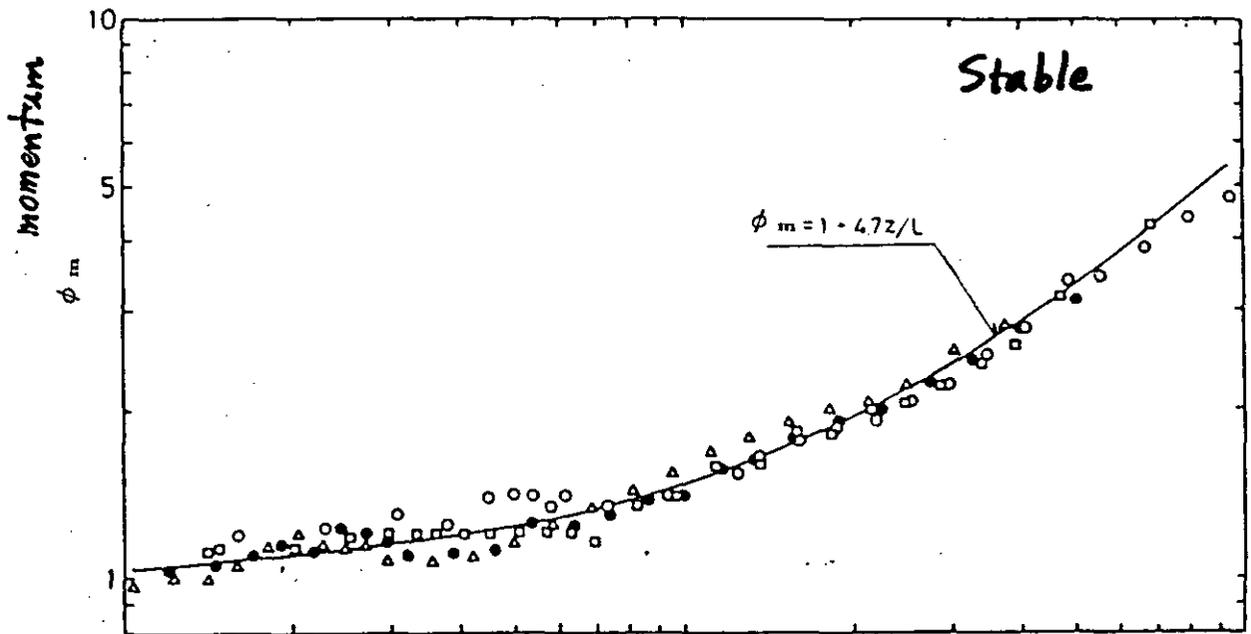
$$\equiv - (u^*)^3 \theta_{ref} / (\kappa g \overbrace{\langle w \theta \rangle}^{u_* T_*})$$

Neutral stratification (adiabatic)

$$z/L = 0 \quad (L \rightarrow \pm \infty)$$



Δ , $Re = 8800$, $\overline{Ri} = -0.0126$; \bullet , $Re = 9840$, $\overline{Ri} = -0.0146$;
 \circ , $Re = 7700$, $\overline{Ri} = -0.0241$.



Δ , $Re = 6600$, $\overline{Ri} = 0.0122$; \bullet , $Re = 6100$, $\overline{Ri} = 0.0140$;
 \square , $Re = 6900$, $\overline{Ri} = 0.0270$; \circ , $Re = 5500$, $\overline{Ri} = 0.0514$.

$$\phi_m \equiv \frac{\kappa z}{u_*} \frac{dU}{dz}, \quad \phi_n = \frac{\kappa' z}{T_*} \frac{d\theta}{dz} = \frac{\kappa' z}{C_*} \frac{dC}{dz} \quad (\kappa' = 1.2\kappa)$$

Integration

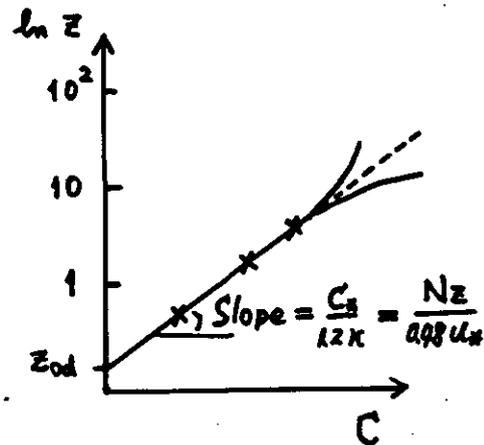
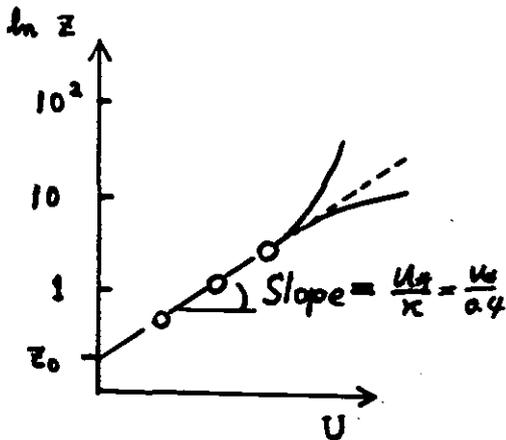
$$\frac{U}{u_*} = \frac{1}{\kappa z} \int_{z_0}^z \frac{1}{\kappa z} \phi_m dz, \quad \frac{C - C_0}{C_*} = \int_{z_{0d}}^z \frac{1}{\kappa' z} \phi_n dz$$

When $\phi_m =$ $\phi_n = 1 + \alpha(z/L)$

$$U/u_* = 1/\kappa [\ln z/z_0 + \alpha(z-z_0)/L]$$

$$(C - C_0)/C_* = 1/(1.2\kappa) [\ln z/z_{0d} + \alpha(z-z_{0d})/L]$$

$$(\theta - \theta_0)/T_* = 1/(1.2\kappa) [\ln z/z_{0h} + \alpha(z-z_{0h})/L]$$



$$z \rightarrow 0 \quad \text{or} \quad L \rightarrow \pm \infty$$

$$U/u_* = \frac{1}{\kappa} \ln \frac{z}{z_0}, \quad \frac{C - C_0}{C_*} = \frac{1}{1.2\kappa} \ln \frac{z}{z_{0d}}$$

Eddy correlation method (covariance method)

$$N_z = -\langle wc \rangle = -\frac{1}{T} \int_0^T wc \, dt = -\frac{1}{N} \sum_{i=1}^N (\tilde{W}_i - \bar{W})(\tilde{C}_i - \bar{C})$$

T : averaging time

\tilde{W}_i, \tilde{C}_i : instantaneous values (digitized)

$(\bar{W}, \bar{C}) = (1/N) \sum (\tilde{W}_i, \tilde{C}_i)$ mean

Instruments: high time resolution (frequency response)
high spatial resolution
rain, snow etc.

Vertical velocity component

Sonic anemometer thermometer ($f_m = 10-20 \text{ Hz}$)

Vane anemometer ($f_m \sim 10 \text{ Hz}$)

Hot-wire anemometer ($f_m = 10^3 \sim 10^4 \text{ Hz}$)

Temperature

Sonic anemometer thermometer

Platinum-wire resistance thermometer

Thermo couple, Quartz thermometer

Water vapor

Lyman- α hygrometer

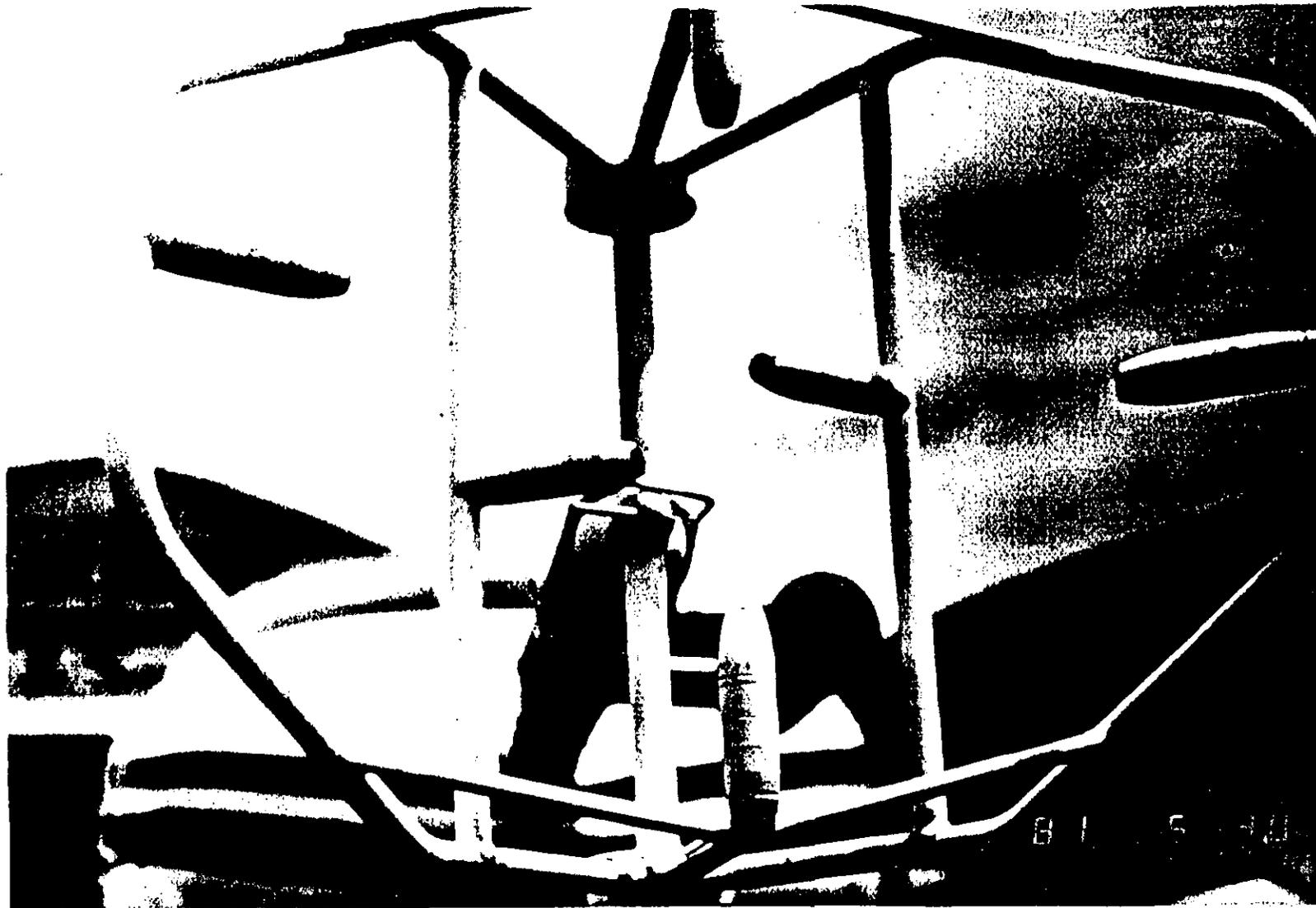
$\text{CO}_2, \text{O}_3, (\text{CH}_3)$

Ultra-red absorption



$$\text{Velocity} = (\text{Prop. time I} - \text{Prop. time II}) / 2L \quad \because \text{Prop. velocity of sound} = \text{Sound vel.} + \text{Wind vel. in still air}$$

$$\text{Temperature} \propto (\text{Prop. time I} + \text{Prop. time II}) / 2L$$



Averaging time \bar{T} (Wyngaard 1973)

$$\bar{T} \approx 2 \tau_i \overline{f'^2} / a^2 \bar{f}^2 \quad \text{Lumley & Panofsky (1964)}$$

\bar{T} : averaging time required to determine the mean f
to an accuracy a

$\overline{f'^2}$: variance of f about its mean

τ_i : integral scale of time function $f \approx \frac{\Lambda}{U} \approx \frac{z}{U}$

$$\bar{T}_{\overline{uw}} \approx \frac{2z}{a^2 U} \left[\frac{\overline{(uw)^2} - (\overline{uw})^2}{(\overline{uw})^2} \right] = \frac{z}{a^2 U} \left[\frac{\overline{(uw)^2}}{u_*^2} - 1 \right]$$

$$\bar{T}_{\overline{wc}} \approx \frac{2z}{a^2 U} \left[\frac{\overline{(wc)^2}}{u_*^2 c_*^2} - 1 \right]$$

$$\bar{T}_{\overline{w^2}} \approx \frac{2z}{a^2 U} \left[\frac{\overline{w^4}}{(\overline{w^2})^2} - 1 \right] \quad \left(\text{Flatness factor } F_w = \frac{\overline{w^4}}{(\overline{w^2})^2} \sim 3 \right)$$

Gaussian.

Neutral

Kansas exp. (tower: 20m)

$$\bar{T}_{\overline{w^2}, \overline{c^2}} \approx 4 \frac{z}{a^2 U} \quad 2.4 \text{ min}$$

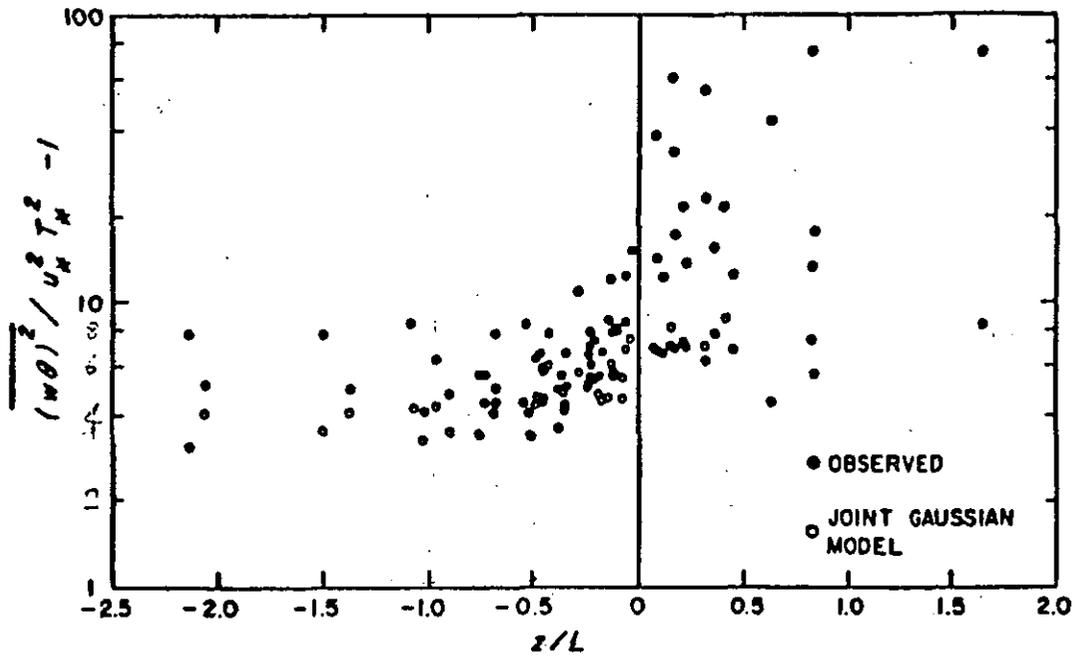
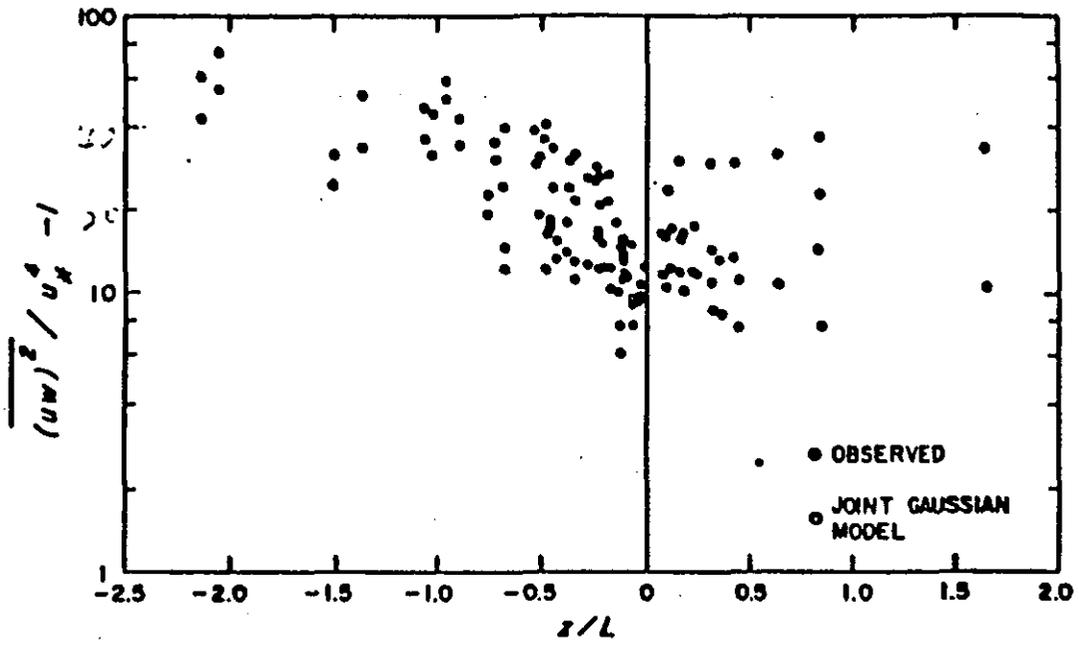
$$\bar{T}_{\overline{wc}, \overline{uw}} \approx 20 \frac{z}{a^2 U} \quad 12 \text{ min}$$

Unstable ($z/L = -1$)

$$\bar{T}_{\overline{w^2}, \overline{c^2}} \approx 4 \frac{z}{a^2 U} \quad 2.4 \text{ min}$$

$$\bar{T}_{\overline{uw}} \approx 100 \frac{z}{a^2 U} \quad 60 \text{ min}$$

$$\bar{T}_{\overline{wc}} \approx 12 \frac{z}{a^2 U} \quad 7.2 \text{ min}$$



Indirect flux measurements

Concentration variance dissipation rate χ_c
(destruction rate of concentration fluctuation)
by molecular diffusion

$$\chi_c = 6 D \left\langle \left(\frac{\partial c}{\partial x} \right)^2 \right\rangle$$

Budget equation for scalar variance

$$-2 \langle w c \rangle \frac{dC}{dz} - \underbrace{\frac{\partial}{\partial z} \langle w c^2 \rangle}_{\text{negligible}} = \chi_c$$

Determination of χ_c

1. $\chi_c \equiv 6 D \left\langle \left(\frac{\partial c}{\partial x} \right)^2 \right\rangle = \frac{6 D}{U^2} \left\langle \left(\frac{\partial c}{\partial t} \right)^2 \right\rangle$ (frozen turbulence)

2. One dimensional spectrum

$$\phi_{cc}(k_1) = \beta_c \chi_c \varepsilon^{-4/3} k_1^{-5/3}, \quad \langle c^2 \rangle = \int_0^{\infty} \phi_{cc}(k_1) dk_1$$

$\beta_c = 0.40$

3. Structure function

$$\langle [C(x+r) - C(x)]^2 \rangle = C_d r^{2/3}$$

In the surface layer (Monin-Obukhov similarity theory)

$$\chi_c = \frac{u_* C_d}{\kappa z} \phi_c \quad \varepsilon = \frac{u_*^3}{\kappa z} \phi_\varepsilon, \quad \chi_o = \frac{u_* T_s}{\kappa z} \phi_o$$

$\phi_c, \phi_o, \phi_\varepsilon$: universal function of z/L

Above the surface layer

Measurements of χ_c and $\frac{dC}{dz}$

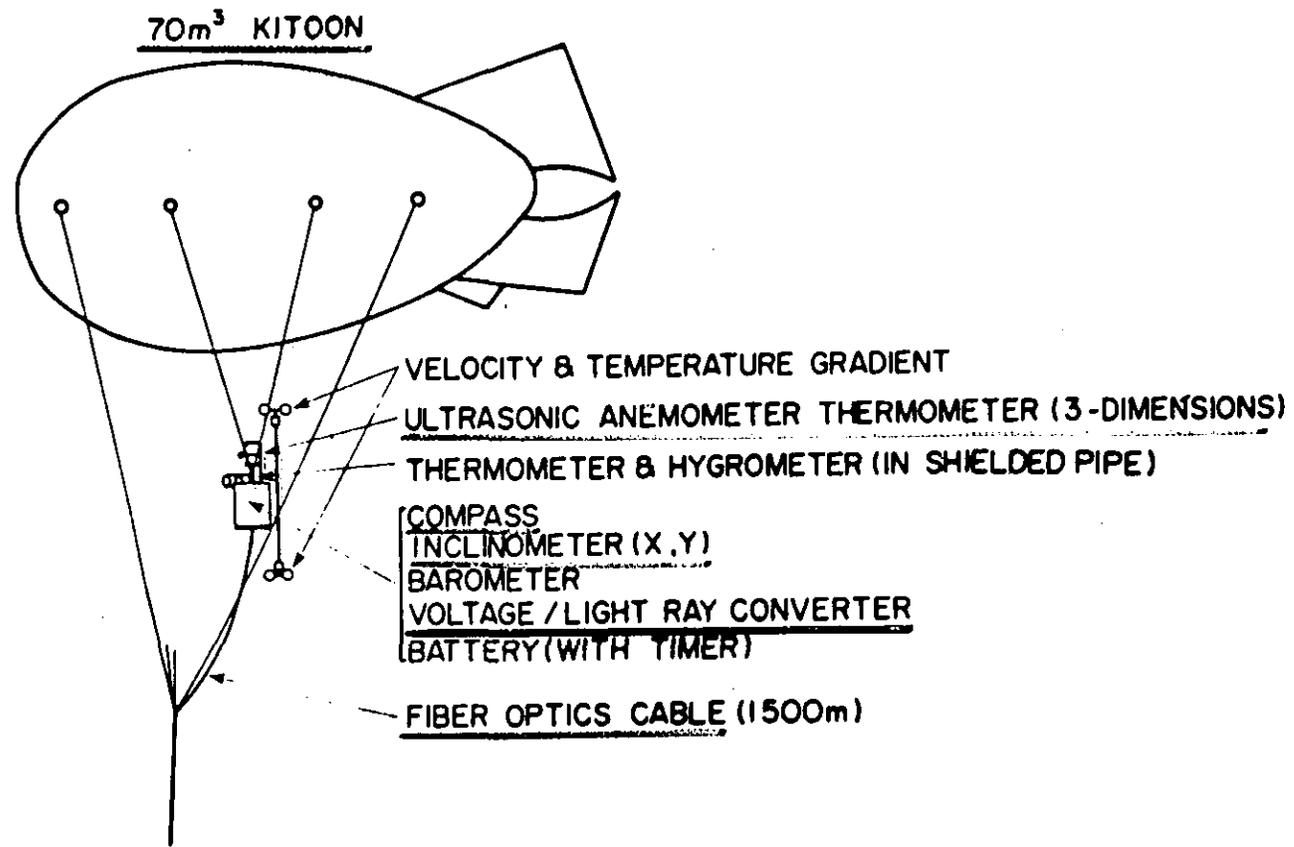
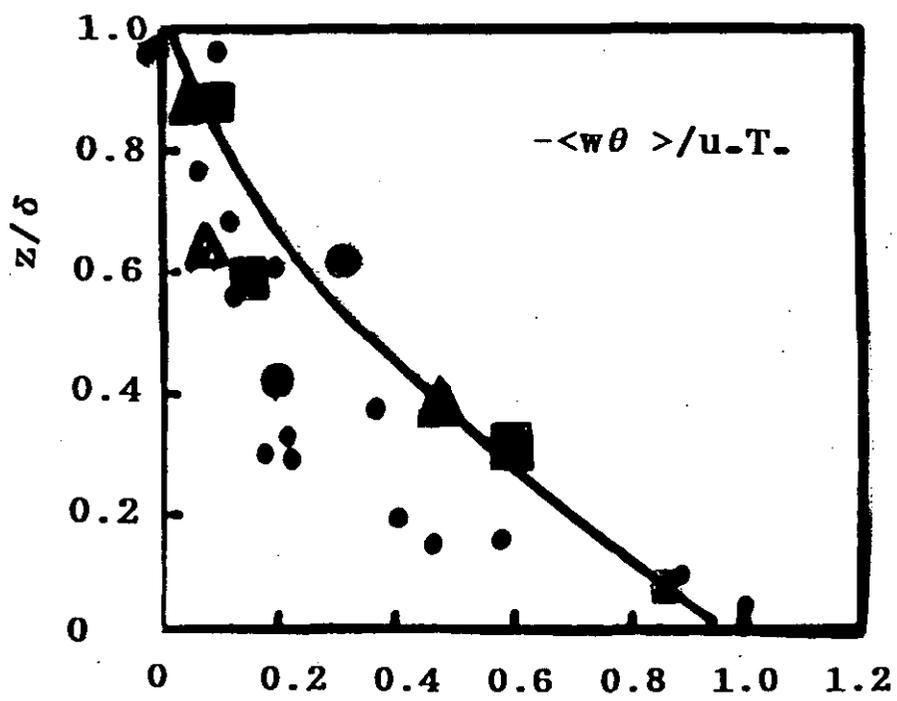
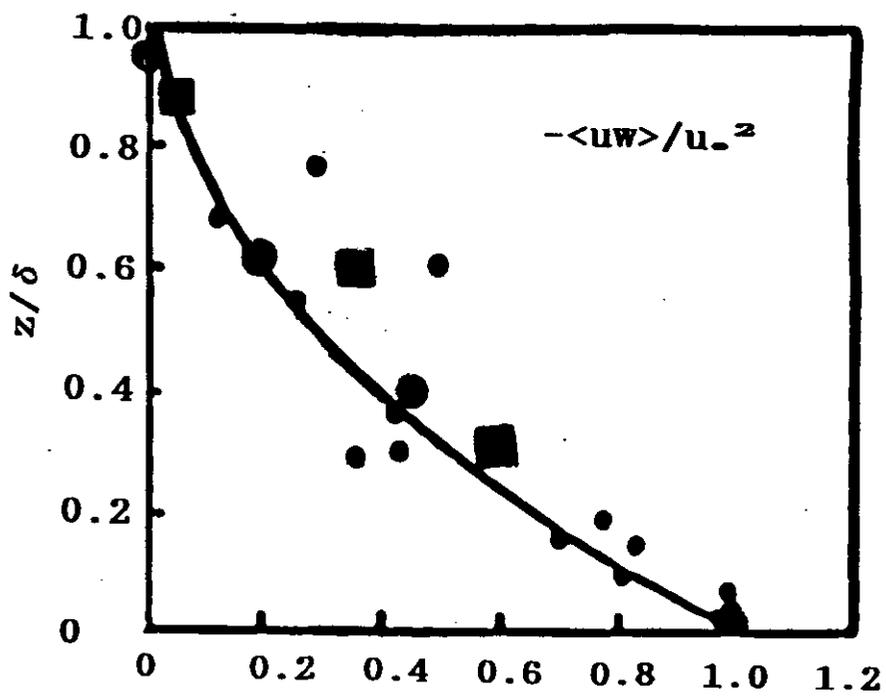


Fig. 1. Schematic of the system.





FACE SUPERIEURE.



FACE INFERIEURE. THIS SIDE UP.



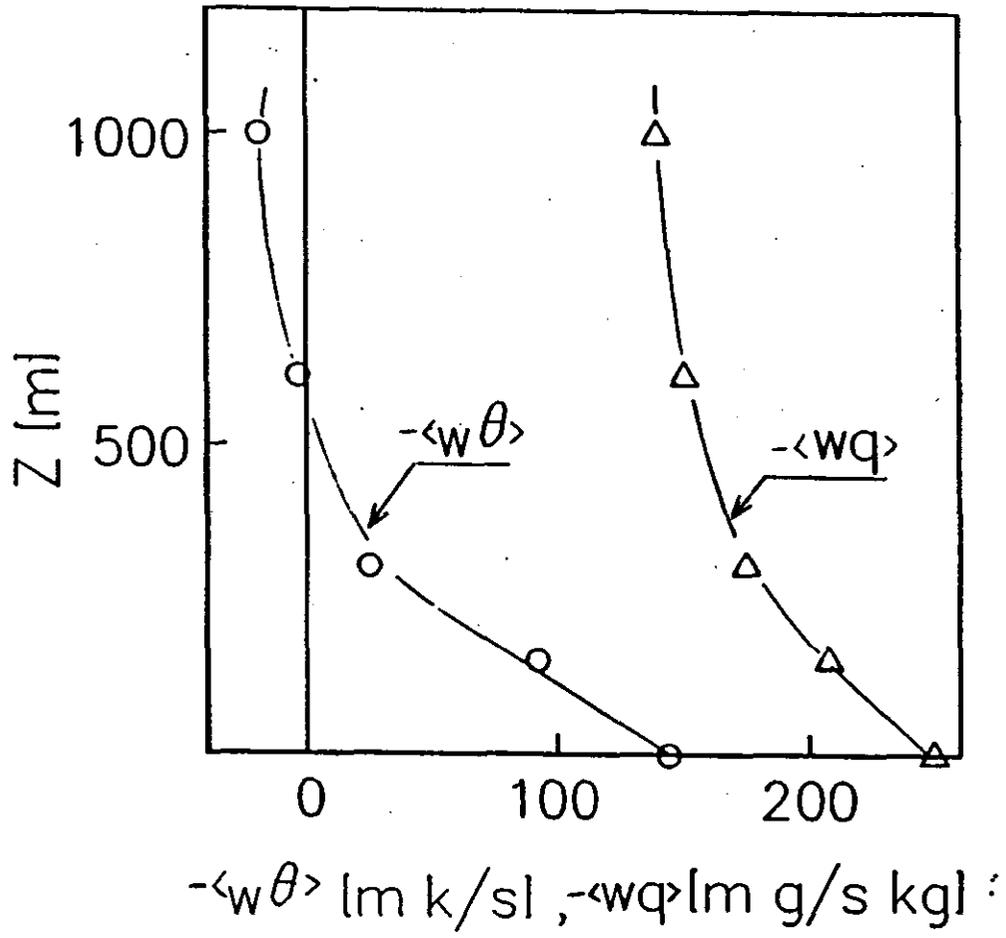
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8-1-5 30

$$\text{Wind velocity against earth} = \text{Air velocity against plane SAT} - \text{Airplane velocity against earth INS} \rightarrow \begin{matrix} u & \text{W-E comp.} \\ v & \text{N-S comp.} \\ w & \text{Vertical comp.} \end{matrix}$$

Channel Contents	Measuring Device	Digital/Analog	Resolution(Nominal)
1 Longitude	INS	Digital	$45^\circ \times 2^{-10}$ (0.309'') (≈ 9.5 m for Longitude)
2 Latitude			
3 Heading Angle			
4 Track Angle			
5 Ground Speed	SAT	Analog	0.1×2^{-6} KNT (≈ 0.08 cm/s)
6 N-S Velocity			
7 E-W Velocity			
8 Pitch	Barometer		0.02°
9 Roll			
10 Vertical Accel.	Radio-altimeter		0.613 cm/s ²
11 X			
12 Y-Z shift	Quartz Thermometer		0.01 m/s
13 Z			
14 Temperature	IRT		0.04° C
15 Altitude			
16 Altitude			0.1 mb
17 Temperature			
18 Ground/Sea Surface Temperature			0.66 m
⋮			0.01° C
⋮			



GAS FLUX MEASUREMENTS BY EDDY CORRELATION

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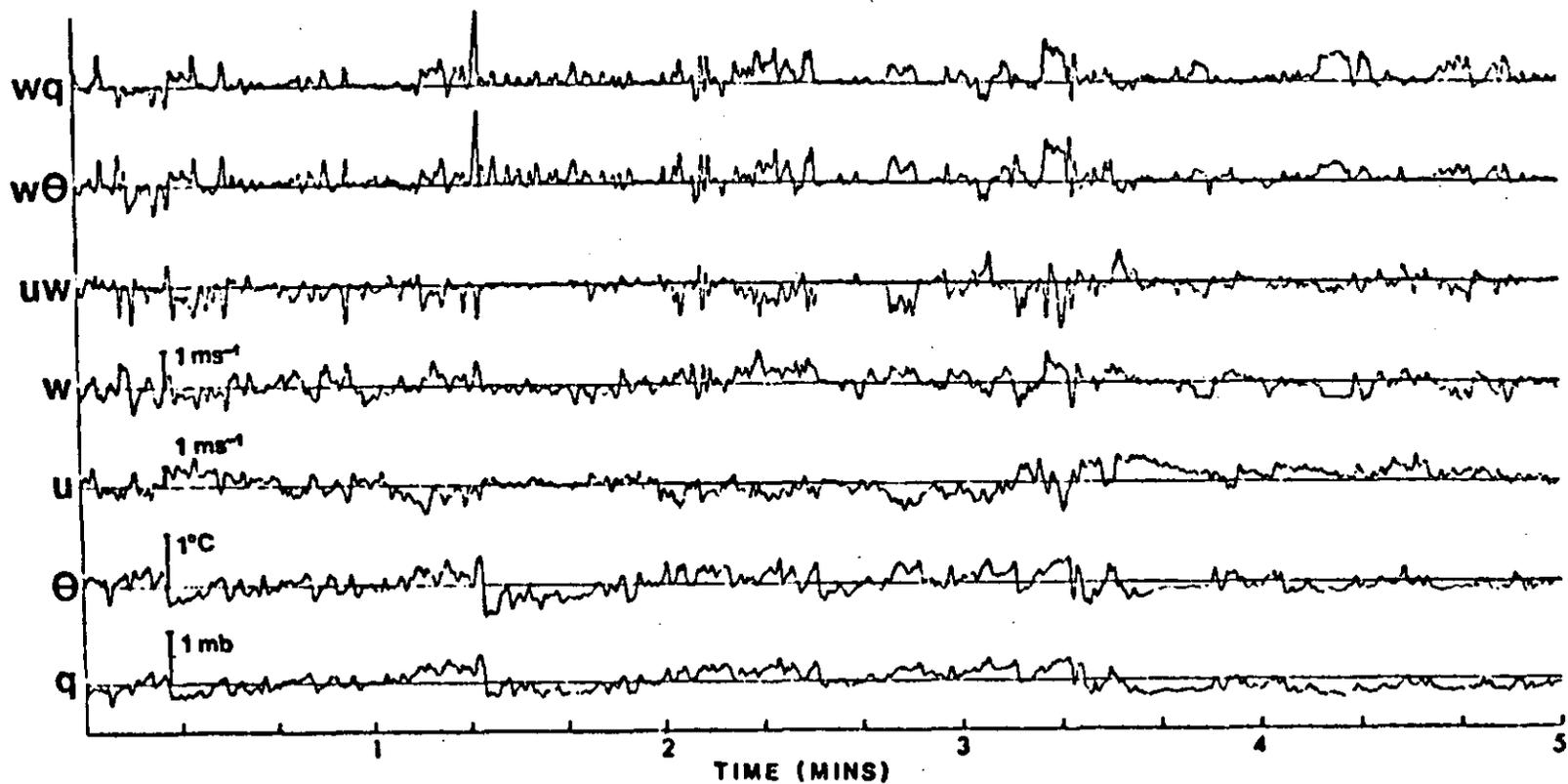
This talk describes some of the issues and problems involved in the field measurement of vertical scalar fluxes, especially trace gas fluxes, by the eddy correlation method. The talk is divided into 6 sections.

1. **Introduction and principle** (overheads 1 to 2): Recalling the basic principle of the eddy correlation method, that a vertical scalar flux F_C in a turbulent flow (with negligible molecular transport compared with turbulent transport) is given by $F_C = \overline{w\rho_c} = \overline{w} \overline{\rho_c} + \overline{w'\rho_c'}$ (where w is the vertical velocity, ρ_c the scalar concentration, overbars denote averages and primes fluctuations).
2. **Bandwidth** (overhead 3): Specifying the frequency range contributing to the eddy covariance.
3. **Webb-Pearman-Leuning (WPL) effect** (overheads 4, 4a): Correcting for the fact that \overline{w} is nonzero over flat ground in thermally non-neutral conditions.
4. **Sampling through a tube** (overheads 5 to 9): Correcting for the damping of scalar fluctuations when air must be drawn through a tube to sample the scalar concentration.
5. **Corrections for platform motion** (overheads 10 to 12): Correcting for ship or aircraft motion.
6. **Avoiding eddy correlation - inverse methods for inferring fluxes from concentrations** (overheads 13 to 15): Introducing two new methods for inferring trace gas fluxes from concentration measurements: an inverse Lagrangian method for deducing source-sink profiles in plant canopies from concentration profiles and turbulence information, and a convective boundary layer (CBL) budget method for inferring trace gas fluxes at regional (100 km²) scales.

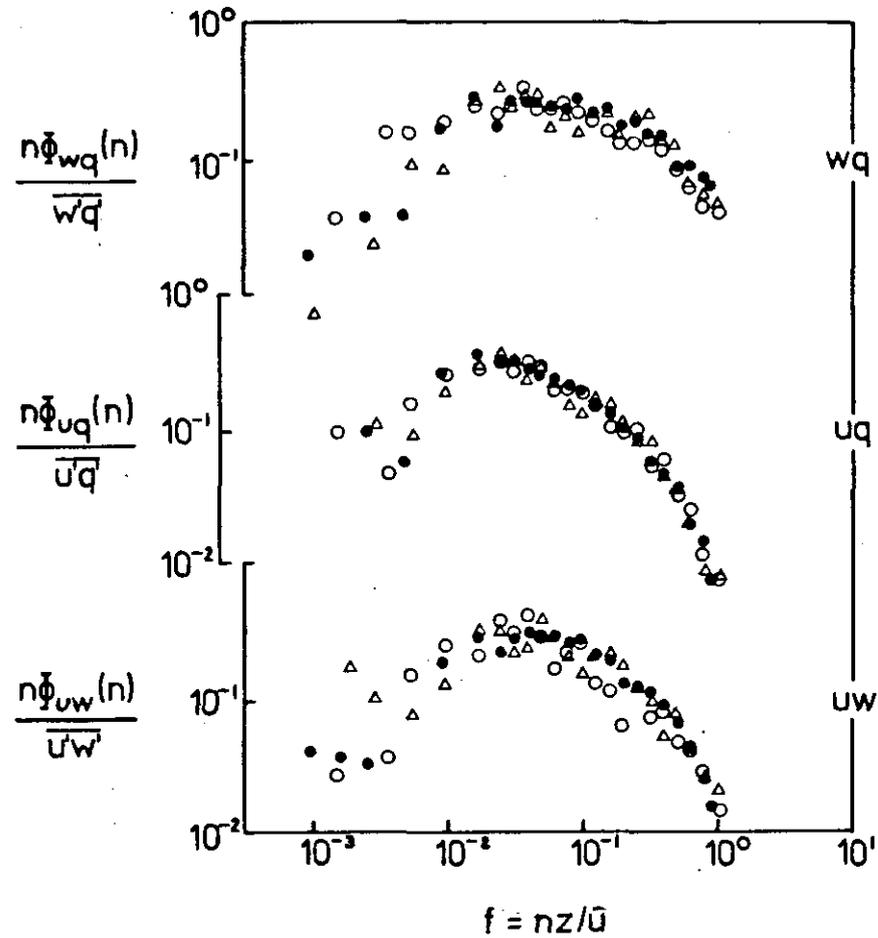
EDDY CORRELATION

1. Principle: $F_c = \overline{w\rho_c} = \overline{w}\overline{\rho_c} + \overline{w'\rho_c'}$
2. Bandwidth: $10^{-3} < n < 10$
 $n = fz/\bar{u}$
 $f = \text{frequency}$
 $z = \text{height}$
 $\bar{u} = \text{mean wind speed}$
3. Webb-Pearman-Leuning effect
4. Sampling along a tube
5. Platform motion
6. Avoiding eddy correlation:
inverse methods for
inferring F_c from $\overline{\rho_c}$

LAKE ALBERT, SOUTH AUSTRALIA, APRIL 1975



LAKE ALBERT, SOUTH AUSTRALIA, 1975



• A △ B ○ C

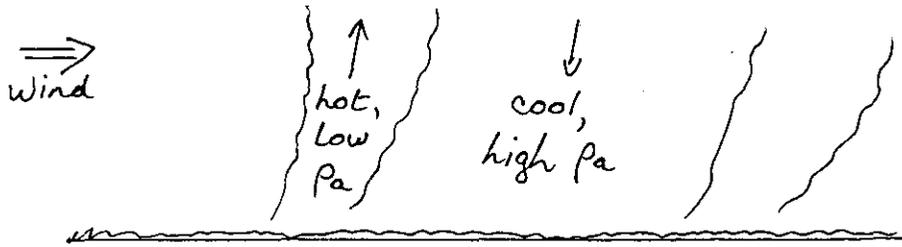
Figure 8. The wq , uq and uw cospectra. The data are grouped in stability bands thus:

band A: $0.11 < -z/L < 0.18$ (4 runs)

band B: $0.18 < -z/L < 0.21$ (3 runs)

band C: $0.27 < -z/L < 0.63$ (5 runs)

WEBB-PEARMAN-LEWING (WPL) EFFECT



Vertical flux of dry air = $\overline{w\rho_a} = 0$

$$\Rightarrow \overline{w\rho_a} + \overline{w'\rho_a'} = 0$$

$$\Rightarrow \overline{w} = -\overline{w'\rho_a'}/\overline{\rho_a}$$

($\overline{w} > 0$ in convection)

For a trace scalar:

$$F_c = \overline{w\rho_c} = \overline{w\rho_c} + \overline{w'\rho_c'}$$

(mass flux) (density) WPL correction

Open path sensor (with temp fluctuations):

$$\overline{w\rho_c} = \overline{\rho_c} \left[\overline{w'm'} + (1 + \overline{m}) \frac{\overline{w'\theta'}}{\overline{T}} \right]$$

Closed path sensor (temp fluctuations removed):

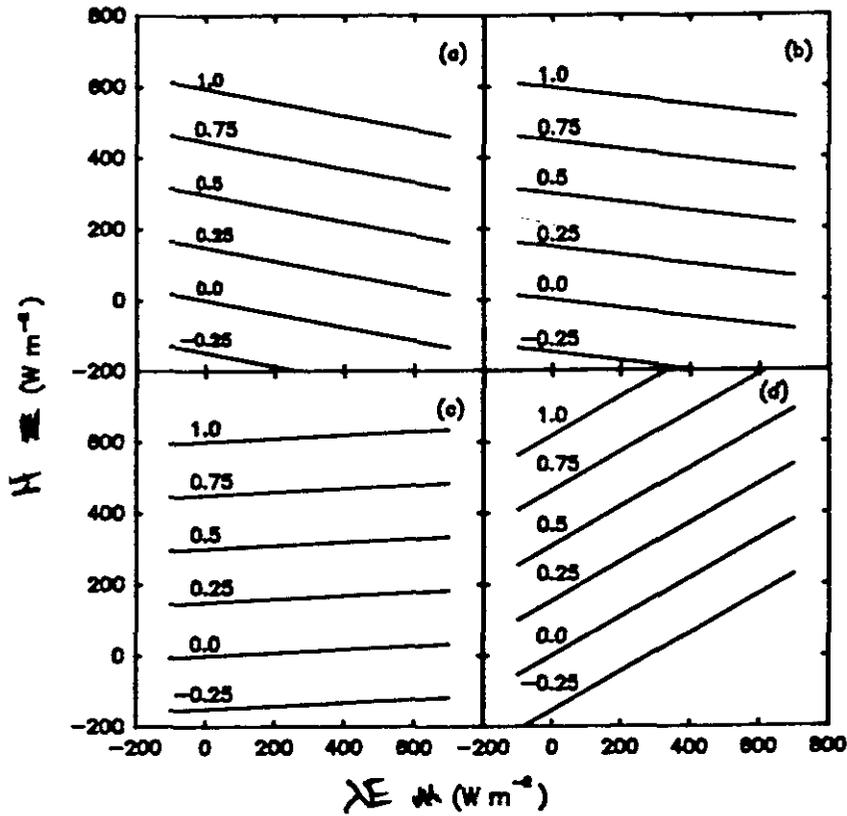
$$\overline{w\rho_c} = \overline{\rho_c} \left[\overline{w'm'} + \overline{m} \frac{\overline{w'\theta'}}{\overline{T}} \right]$$

with $m = \text{H}_2\text{O}$ molar mixing ratio.

$\overline{w\rho_c}/F_c$ range:

-0.1 to ≈ -1	(CO ₂)
0.25 to 10	(N ₂ O)
0.04 to 0.6	(CH ₄)

LEUNING AND MONCRIEFF (1990)



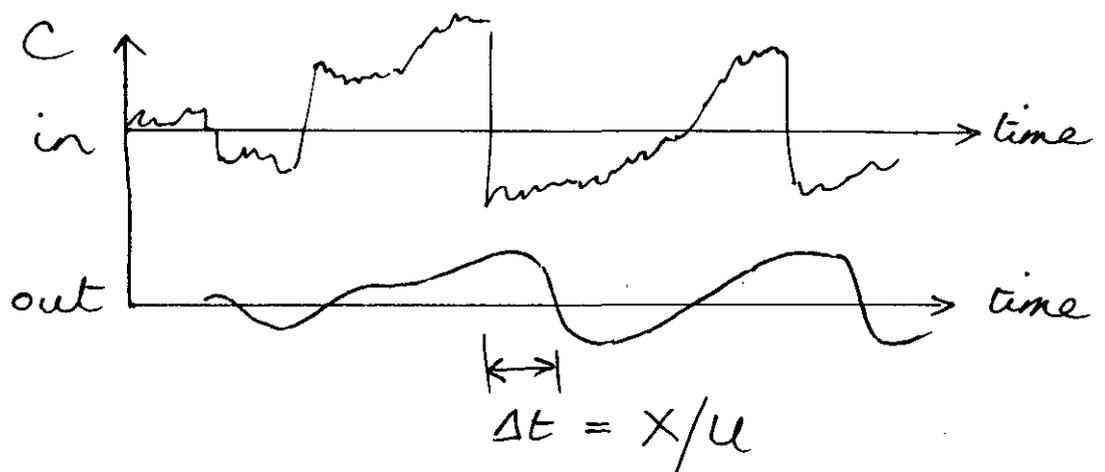
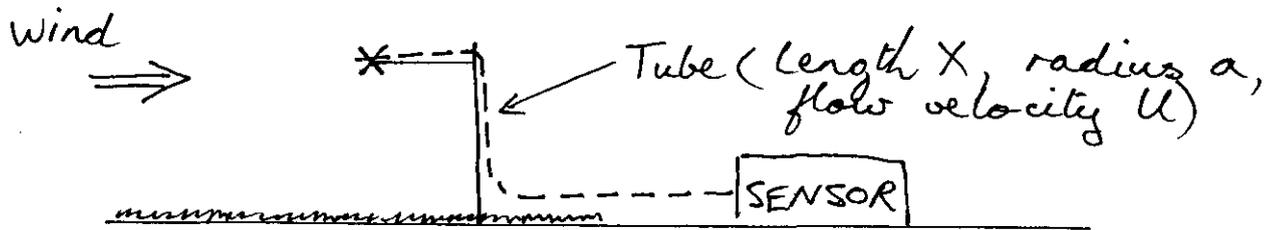
$\bar{w}_p c$ for CO_2
($\text{mg m}^{-2} \text{s}^{-1}$)

a → b → c → d
represents
increasing
 H_2O cross-
sensitivity

Fig. 2. Isopleths for the correction to measured CO_2 fluxes as a function of latent and sensible heat fluxes for open-path CO_2 analysers with differing cross-sensitivity to water vapour. (a) $\beta/\alpha = 0$, (b) $\beta/\alpha = 3 \times 10^{-4}$, (c) $\beta/\alpha = 1 \times 10^{-3}$ and (d) $\beta/\alpha = 3 \times 10^{-3}$. Numbers on the lines represent the flux correction to be added to F_{raw} ($\text{mg CO}_2 \text{ m}^{-2} \text{ s}^{-1}$). Corrections at $H = 0$ also apply to closed-path CO_2 analyser.

SAMPLING THROUGH A TUBE

Lenschow and Raupach, J. Geophys. Res. (1991)
96, 15259-68.



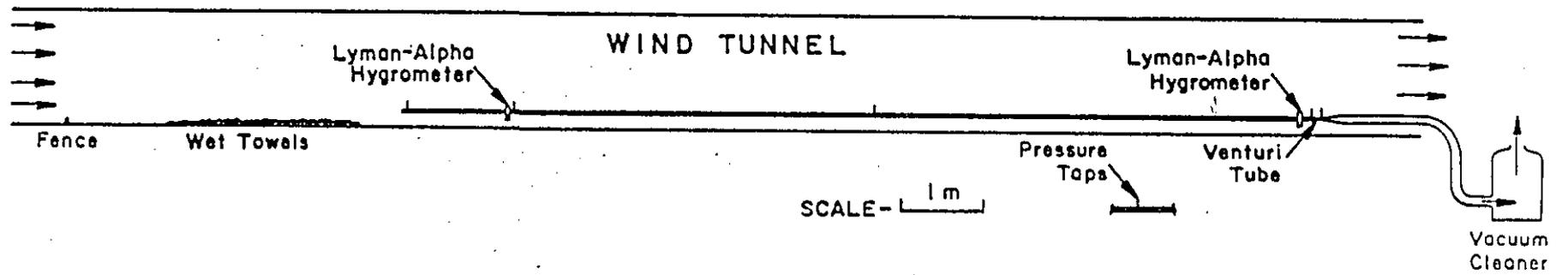
Tube transfer function:

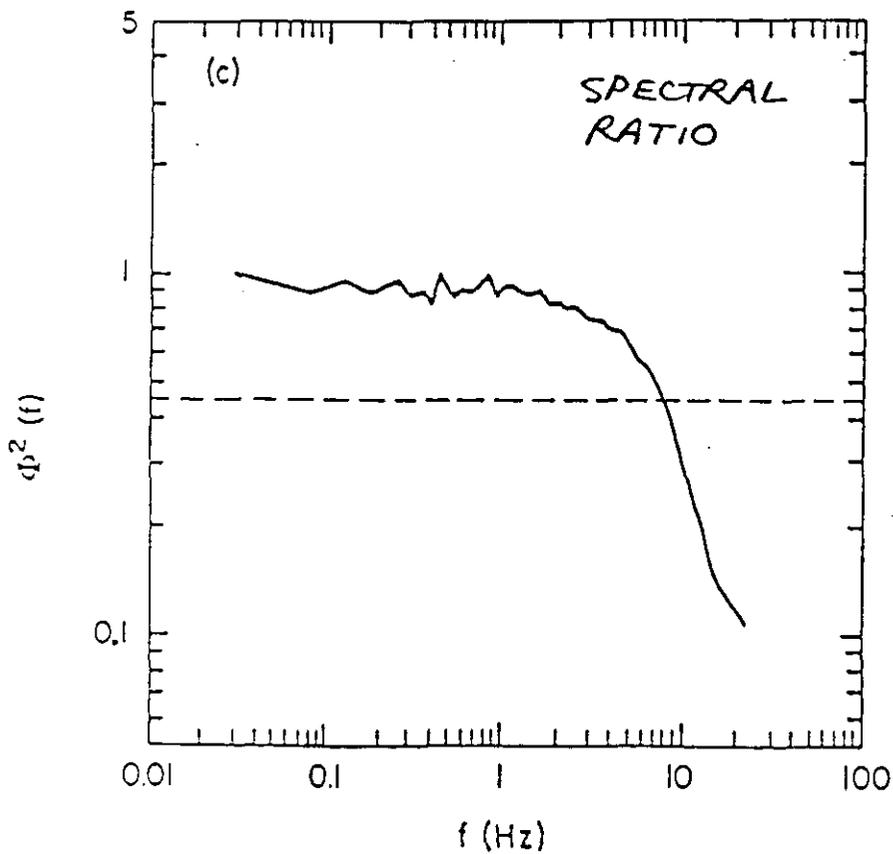
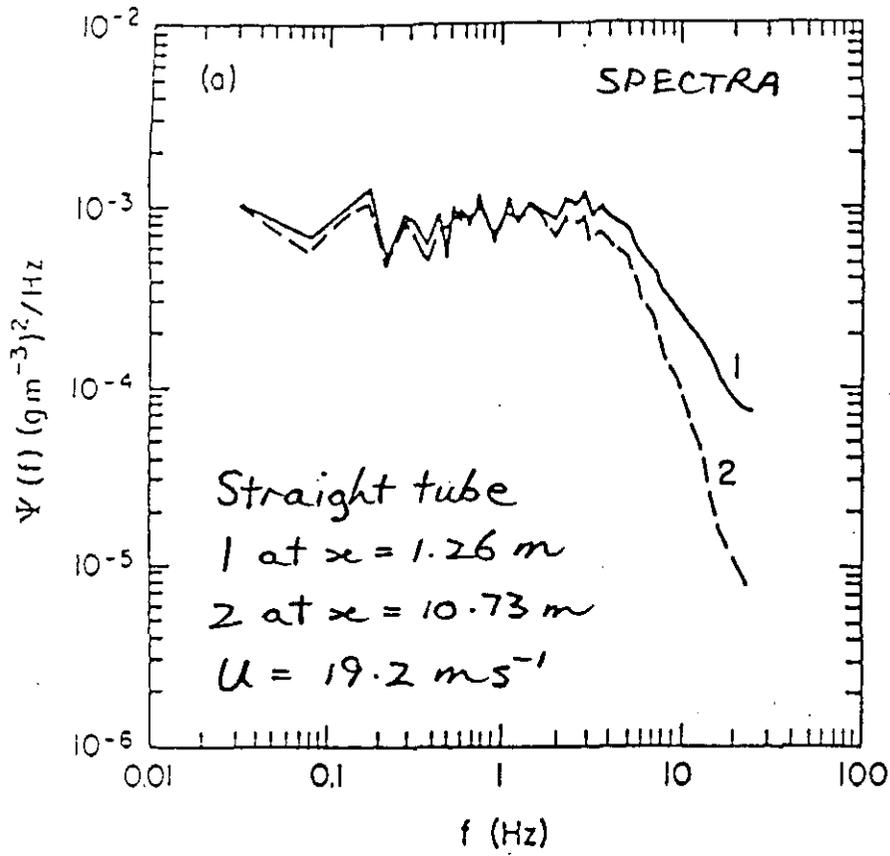
$$\Phi(n) = \frac{\text{amplitude out}}{\text{amplitude in}} = \exp\left(\frac{-n^2 \ln \sqrt{2}}{n_0^2}\right)$$

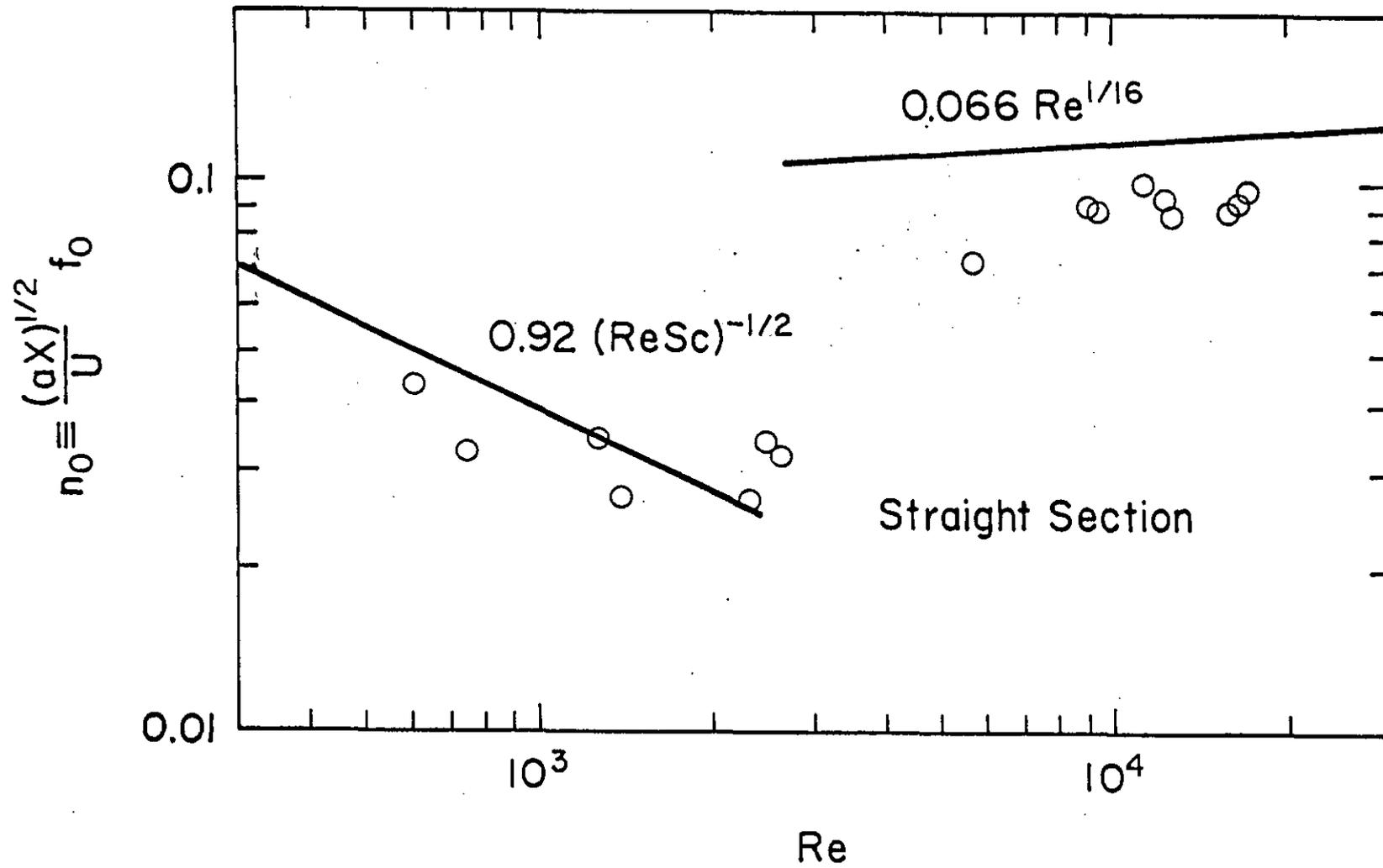
$$n = \frac{\omega}{U} \sqrt{aX} = \text{dimensionless frequency}$$

$$n_0 = \begin{cases} 0.92 (\text{Re Sc})^{-1/2} & (\text{laminar}) \\ 0.066 \text{Re}^{1/16} & (\text{turbulent}) \end{cases}$$

Desirable: turbulent flow in tube!
HIGH measurement point!







LEUNING AND MONCRIEFF (1990)

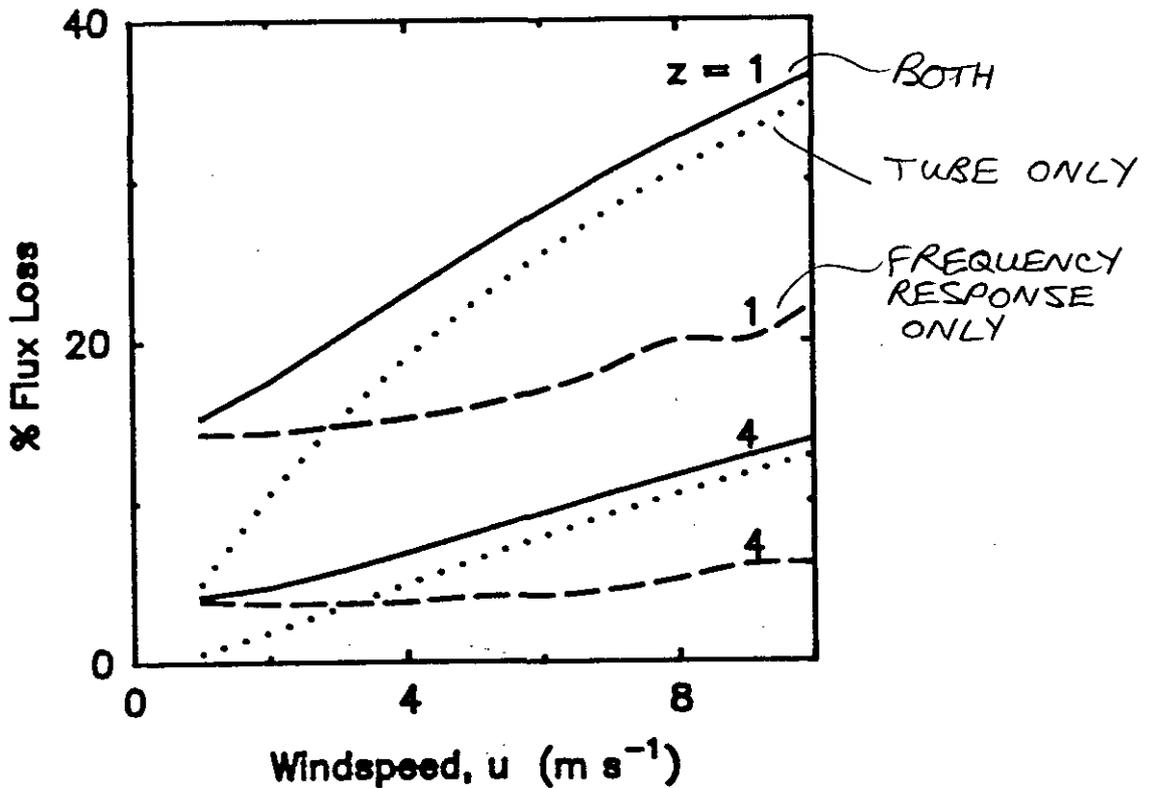
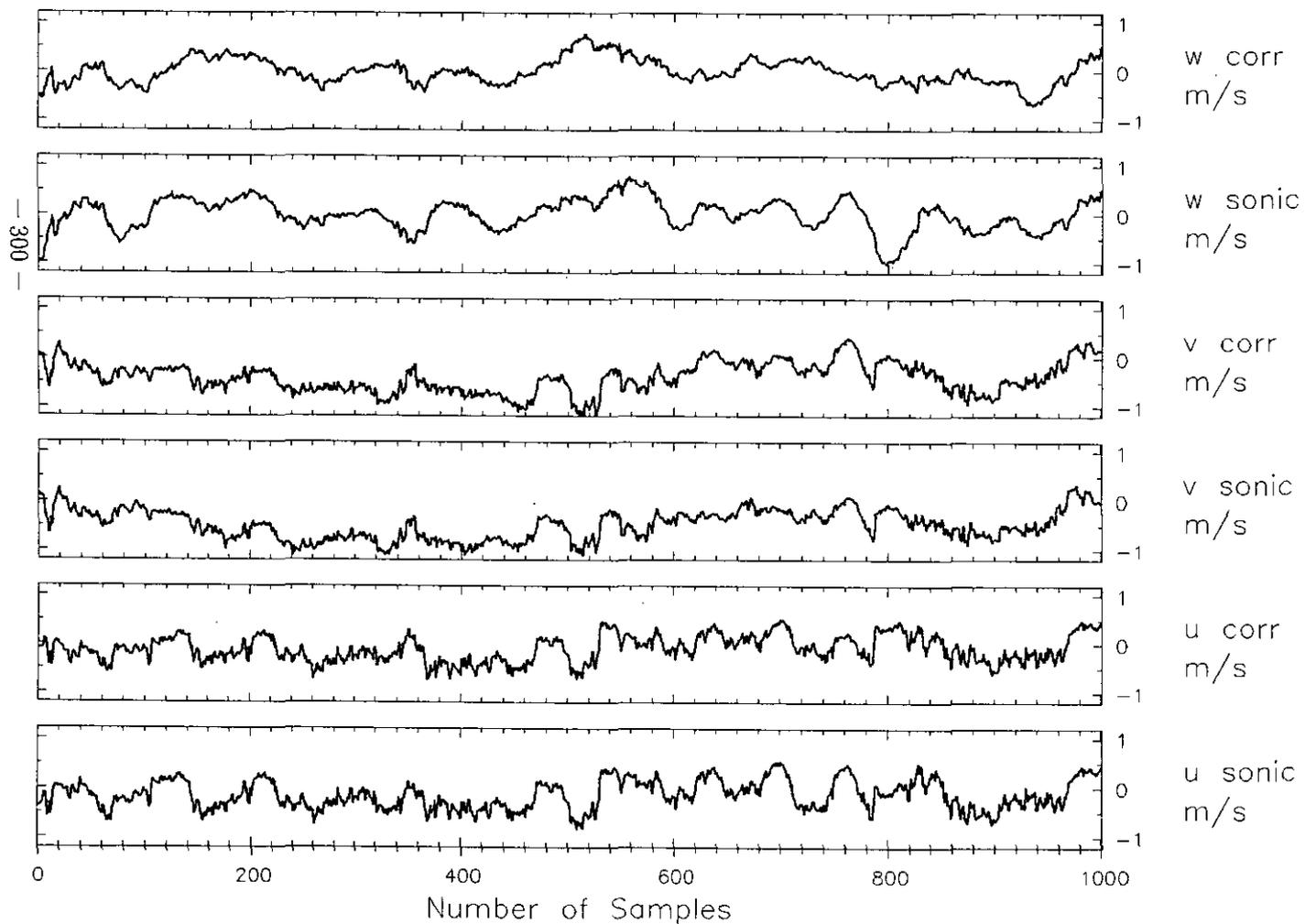
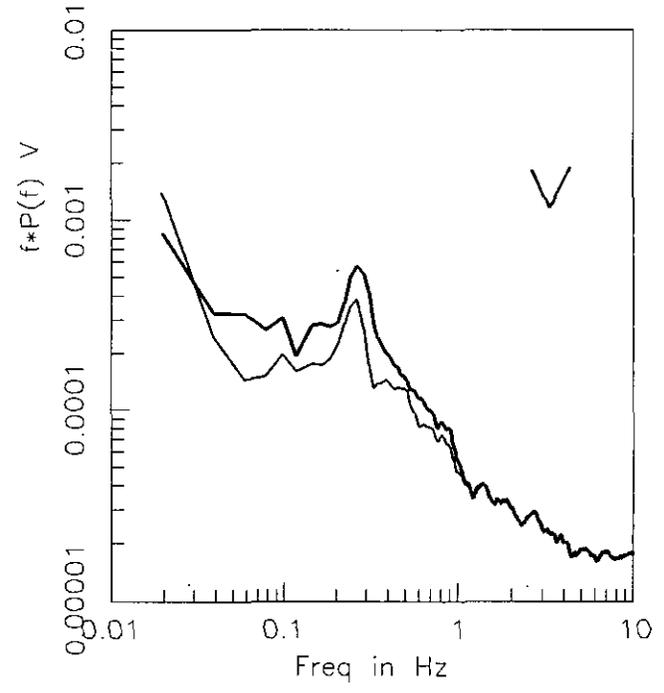
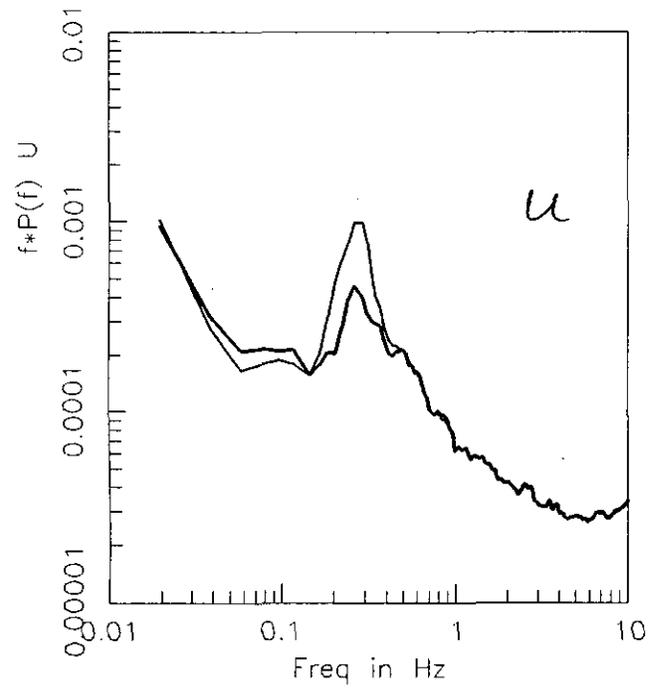
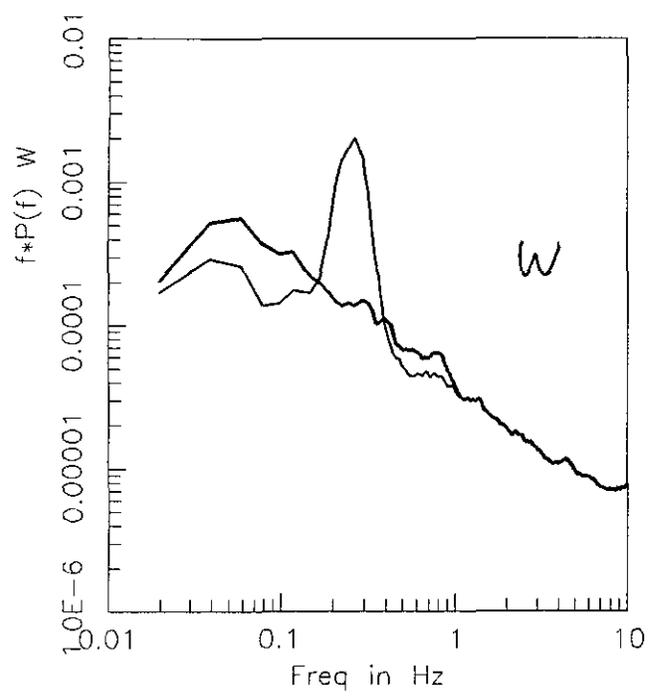


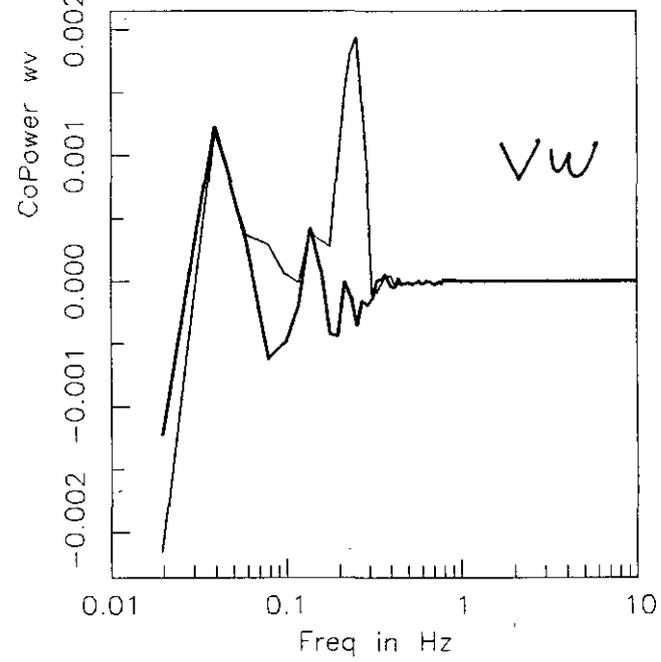
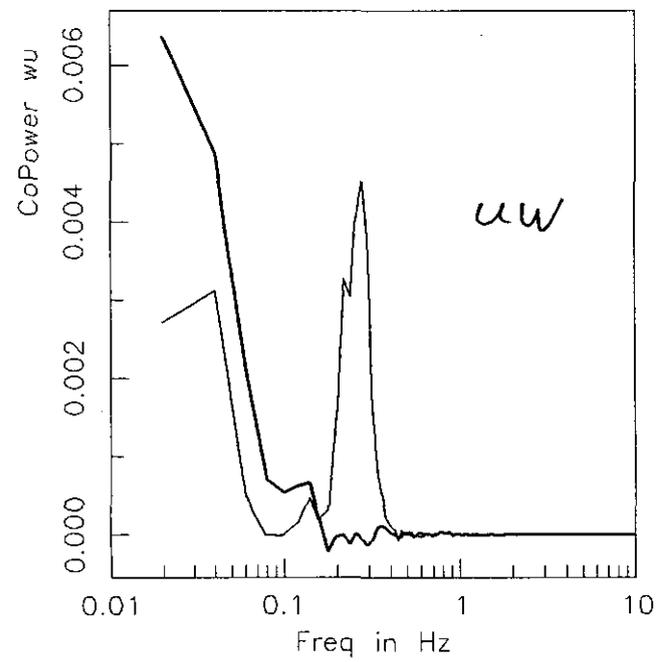
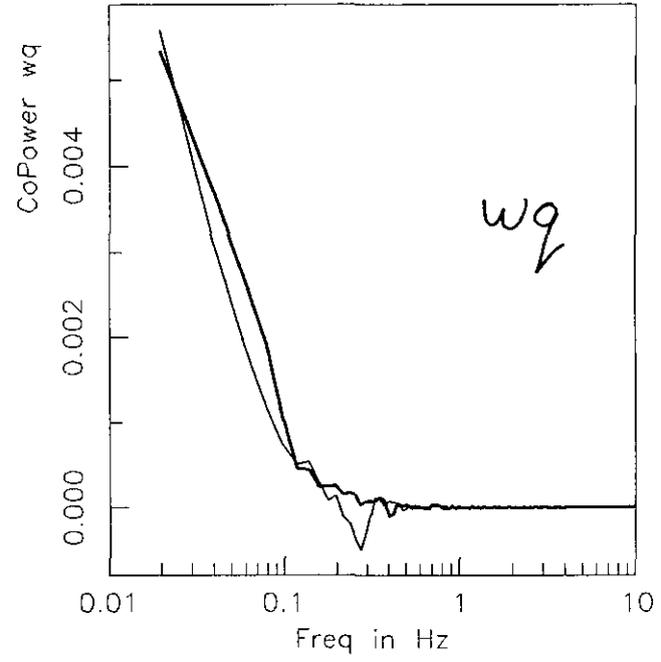
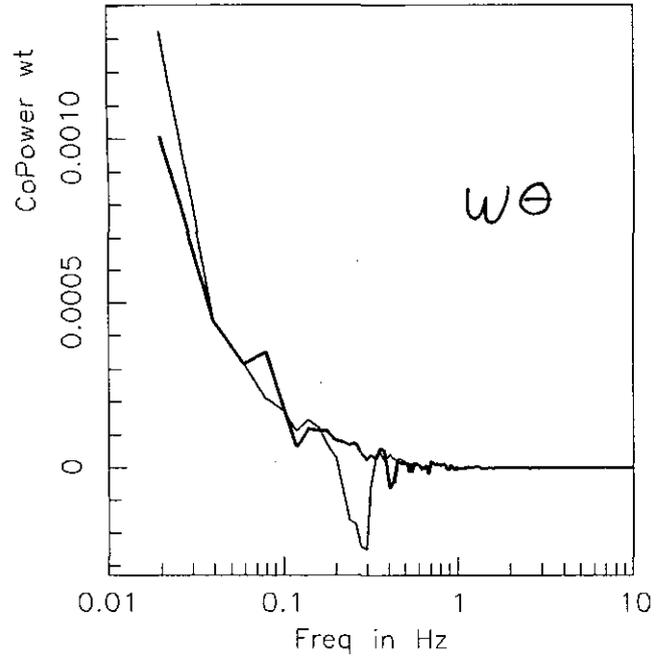
Fig. 3. Estimated percentage loss in measured F_{raw} for a closed-path CO₂ analyser supplied with air sampled through a copper tube $L = 1.5$ m, $r_0 = 3$ mm and $V = 8 \times 10^{-3}$ s⁻¹. Separation distance between tube inlet and sonic anemometer was 0.18 m. Losses are plotted against windspeed for sensor heights of 1 and 4 m above the zero-plane. Flux losses due to damping of turbulence fluctuations by tube only (.....); flux losses due to limited sensor frequency response but neglecting effect of tube (----) and combined influence of tube and sensor frequency response (—).

Correction for ship motion:
Shao, Bradley, Coppin (current work)

Sonic and Corrected Wind: Sonic1-0535







INVERSE METHODS

- inferring fluxes from mean concentrations

1. Canopy scale

$$C_i - C_R = \sum_{j=1}^m D_{ij} S_j \Delta z_j$$

concentration at level i

source strength in layer j

dispersion matrix: depends on σ_w, T_L

2. Scale of planetary convective B.L.

CBL well mixed

$$\frac{dC}{dt} = \frac{F_c}{h} + (C_+ - C) \frac{dh}{dt}$$

$$\Rightarrow \int_0^t F_c dt = (C - C_+) h(t)$$

(with assumptions)

